

Metasurface Synthesis

Using Momentum Transformation

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We introduce a novel approach to model spatial discontinuities, such as thin surfaces and metasurfaces, in the momentum space (spatial spectrum). The approach is called “momentum transformation” to emphasize its spatial spectral nature and avoid confusion with temporal spectral techniques.

The momentum transformation approach is analogous to expressing Maxwell’s equations in the sense of distributions* but in the momentum space and is applicable to scalar and vector waves. The momentum transformation is particularly suitable for the scalar case due to its simplicity.

The momentum transformation has been applied to metasurface synthesis, where it yields a complete description of the metasurface in the momentum space and provides physical insight into the transformation operation.

These slides are adopted from an oral presentation at URSI GASS 2014, Beijing, China.

Keywords: 3D Metamaterials, Metasurfaces, Frequency Selective Surfaces, Metasurface Synthesis Technique, Momentum Conservation, Comparison with Optics, Momentum Transformation.

I. MOTIVATION

II. MT METHOD

III. COMPARISON WITH OPTICS

IV. EXAMPLES

❖ CONCLUSIONS & QUESTIONS

I. MOTIVATION

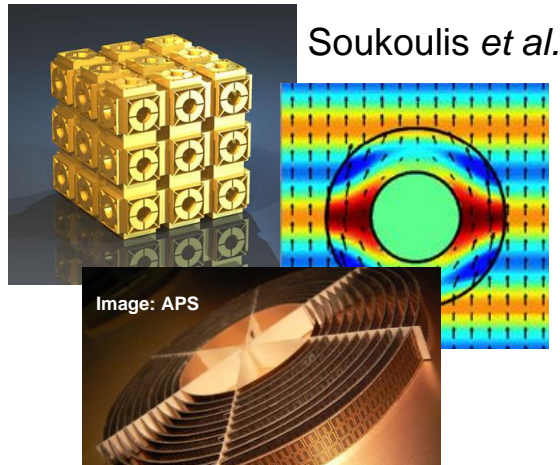
II. MT METHOD

III. COMPARISON WITH OPTICS

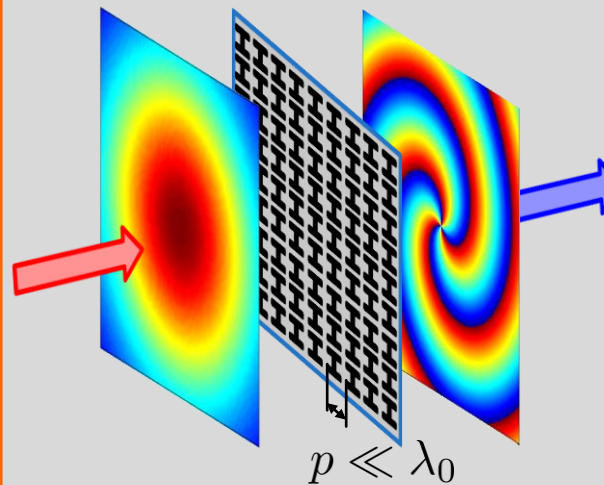
IV. EXAMPLES

❖ CONCLUSIONS & QUESTIONS

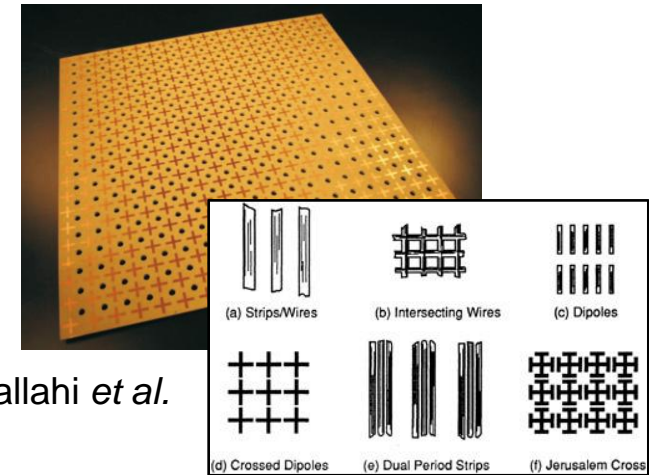
3D Metamaterials



Metasurfaces

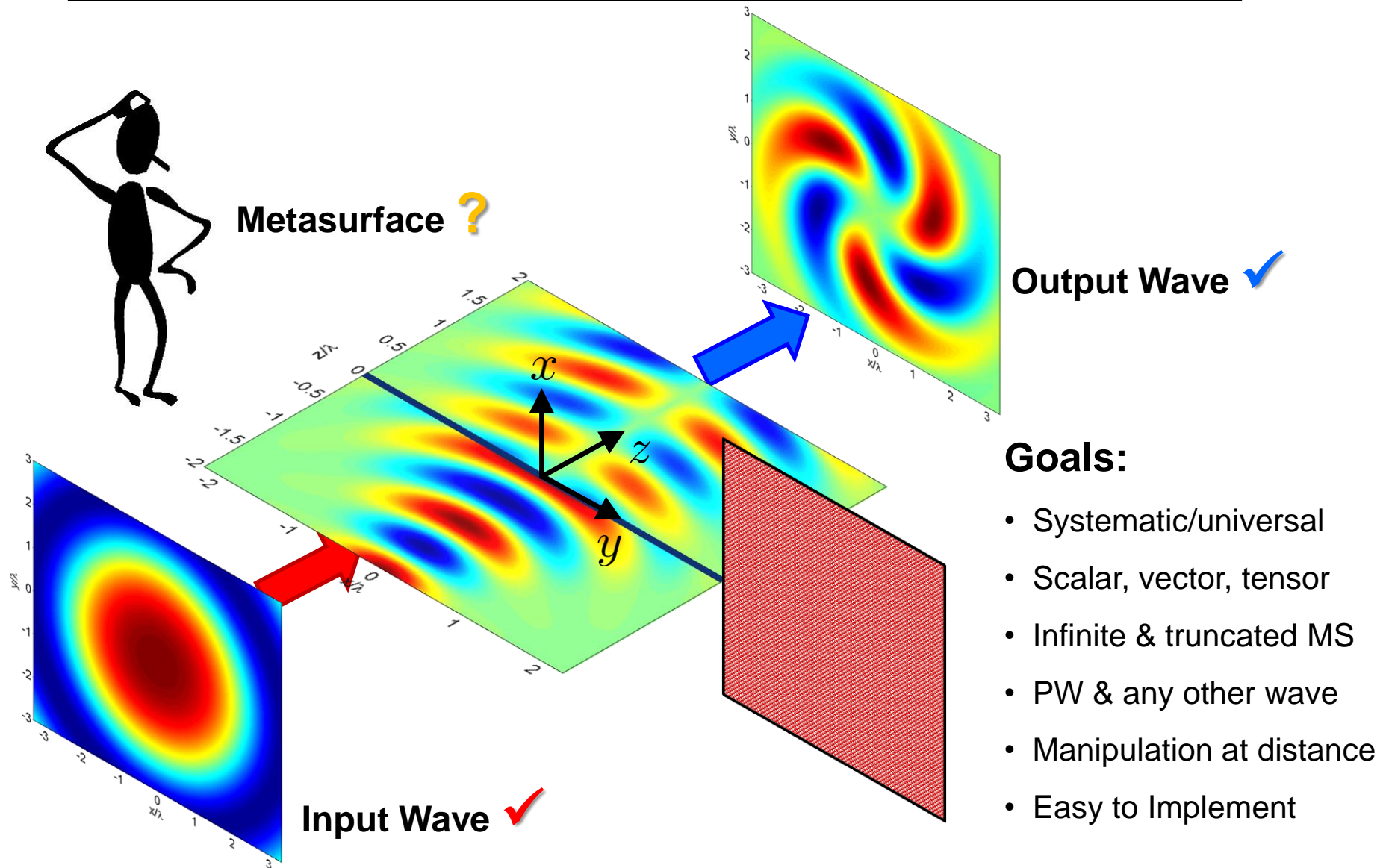


Frequency Selective Surfaces (FSS)



- | | | | | |
|--------------------------|---|--------------------------|---|-----------------------------|
| ☺ Extra Field Control | ← | ☺ Extra Field Control | → | ☺ Reduced Dimensionality |
| ☺ Signal Processing | ← | ☺ Reduced Dimensionality | → | ☺ Filtering & Polarization |
| ☹ Bulkiness | ← | ☺ Extensive Processing | → | ☹ Limited Signal Processing |
| ☹ High Loss | ← | ☺ Low Loss | → | ☹ Uniformity Restriction |
| ☹ Fabrication Difficulty | ← | ☺ Fabrication Easiness | → | |
| | | ☺ Incidence Flexibility | → | |

Metasurface Synthesis Technique ?



Goals:

- Systematic/universal
- Scalar, vector, tensor
- Infinite & truncated MS
- PW & any other wave
- Manipulation at distance
- Easy to Implement

Benefits:

- Physical insight
- Inspiring new devices

SYNTHESIS → TWO-STEPS:

- Determination of metasurface transfer function
- Implementation (scattering particles)

I. MOTIVATION

II. MT* METHOD

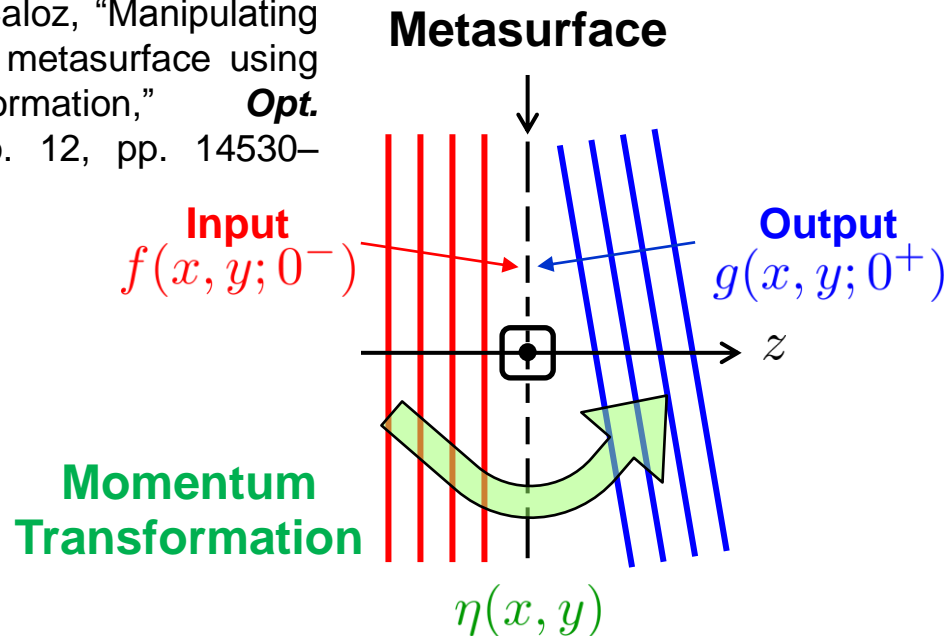
III. COMPARISON WITH OPTICS

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❖ CONCLUSIONS & QUESTIONS

MT*: Momentum Transformation
(k , not “spectral” avoid confusion with ω)

M. A. Salem and C. Caloz, “Manipulating light at distance by a metasurface using momentum transformation,” *Opt. Express*, vol. 22, no. 12, pp. 14530–14543, Jun. 2014.



Linear shift-variance: $g(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x', y') h(x', y'; x, y) dx' dy'$

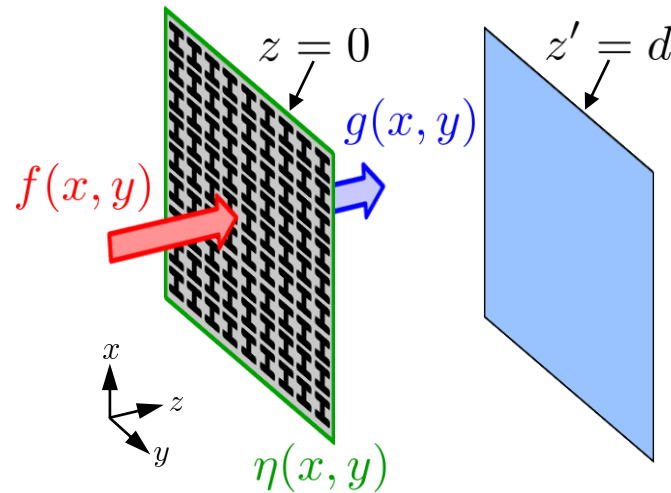
Zero EM thickness \Rightarrow locality: $h(x', y'; x, y) = \delta(x' - x, y' - y) \eta(x', y')$

Consequence of locality: $g(x, y) = f(x, y) \eta(x, y)$

Fourier transform: $\tilde{g}(k_x, k_y) = \tilde{f}(k_x, k_y) * \tilde{\eta}(k_x, k_y)$



Momentum shift-invariance \equiv momentum conservation: $\tilde{f}(k_x - K_x, k_y - K_y) \xrightarrow{z=0} \tilde{g}(k_x - K_x, k_y - K_y)$



Locality: $g(x, y) = \eta(x, y) f(x, y) \rightarrow \eta(x, y) = \frac{g(x, y)}{f(x, y)}$

Metasurface: $\eta(x, y) = g(x, y) \xi(x, y), \quad \xi(x, y) = \frac{1}{f(x, y)}$

Momentum conservation: $\tilde{\eta}(k_x, k_y) = \tilde{g}(k_x, k_y) * \tilde{\xi}(k_x, k_y)$

Plane wave (Fourier) expansion at $z = 0$ of $\psi = f, \eta, g$:

$$\psi(x, y) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \tilde{\psi}(k_x, k_y) e^{+i(k_x x + k_y y)} dk_x dk_y$$

$$\tilde{\psi}(k_x, k_y) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \psi(x, y) e^{-i(k_x x + k_y y)} dx dy$$

Momentum wave equation:

$$k^2 = k_x^2 + k_y^2 + k_z^2$$

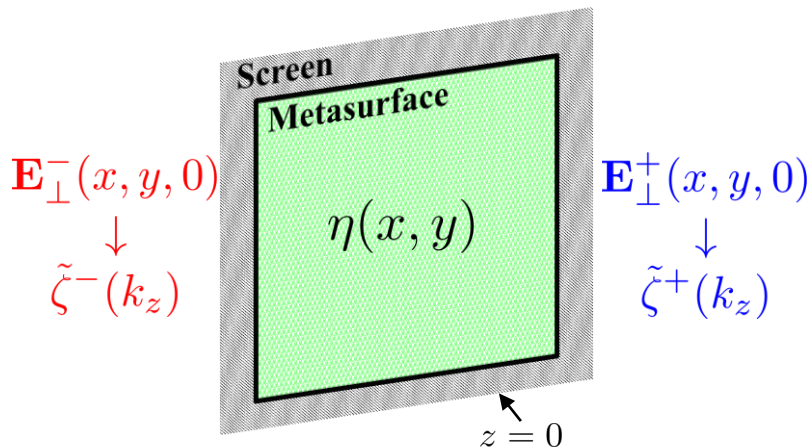
Specification at $z = d$: $g(x, y; z' = d) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \underbrace{\tilde{g}(k_x, k_y) e^{-ik_z d}}_{\tilde{g}_d(k_x, k_y)} e^{+i(k_x x + k_y y)} dk_x dk_y (e^{-i\omega t})$

Reverse

propagation: $\tilde{g}(k_x, k_y) \rightarrow \tilde{g}_d(k_x, k_y) = \tilde{g}(k_x, k_y) \tilde{\varphi}(k_x, k_y; k; d), \quad \tilde{\varphi}(k_x, k_y; k; d) = e^{-i\sqrt{k^2 - k_x^2 - k_y^2} d}$

Momentum transformation equation: $\tilde{\eta}(k_x, k_y) = \tilde{g}(k_x, k_y) \tilde{\varphi}(k_x, k_y; k, d) * \tilde{\xi}(k_x, k_y)$

VPW orthogonality:
$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \mathbf{E}_{\perp}^{\text{VPW}}(x, y; k_z) \times \mathbf{H}_{\perp}^{*\text{VPW}}(x, y; k'_z) dx dy = P_{k_z}^{\text{VPW}} \delta(k_z - k'_z)$$



k_z eigenmode expansion:

$$\mathbf{E}^{\mp}(x, y; z) = \int_0^{+\infty} \tilde{\zeta}^{\mp}(k_z) \mathbf{E}_{\perp}^{\text{VPW}}(x, y; k_z) e^{ik_z z} dk_z$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \dots \times \mathbf{H}^{*\text{VPW}\mp}(x, y; k_z) dx dy$$

& VPW orthogonality

**Scalar
expansion
coefficients:**

$$\tilde{\zeta}^-(k_z) = \frac{1}{P_{k_z}^{\text{VPW}}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \mathbf{E}_{\perp}^-(x, y) \times \mathbf{H}_{\perp}^{*\text{VPW}}(x, y; k_z) dx dy$$

$$\tilde{\zeta}^+(k_z) = \frac{1}{P_{k_z}^{\text{VPW}}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \mathbf{E}_{\perp}^+(x, y) \times \mathbf{H}_{\perp}^{*\text{VPW}}(x, y; k_z) dx dy$$

Reverse propagator: $\tilde{\varphi}(k_x, k_y; k, d) = e^{-ik_z d}$

Scalarized (correspondence with scalar case) solution: $\tilde{\zeta}^-(k_z) * \tilde{\eta}(k_z) = \tilde{\zeta}^+(k_z) \tilde{\varphi}(k_z, d)$

Modal transformation, using $k^2 = k_x^2 + k_y^2 + k_z^2$:
$$\tilde{\eta}(k_x, k_y) = \tilde{\eta}(k_z) = \sqrt{k^2 - k_x^2 - k_y^2}$$

I. MOTIVATION

II. MT METHOD

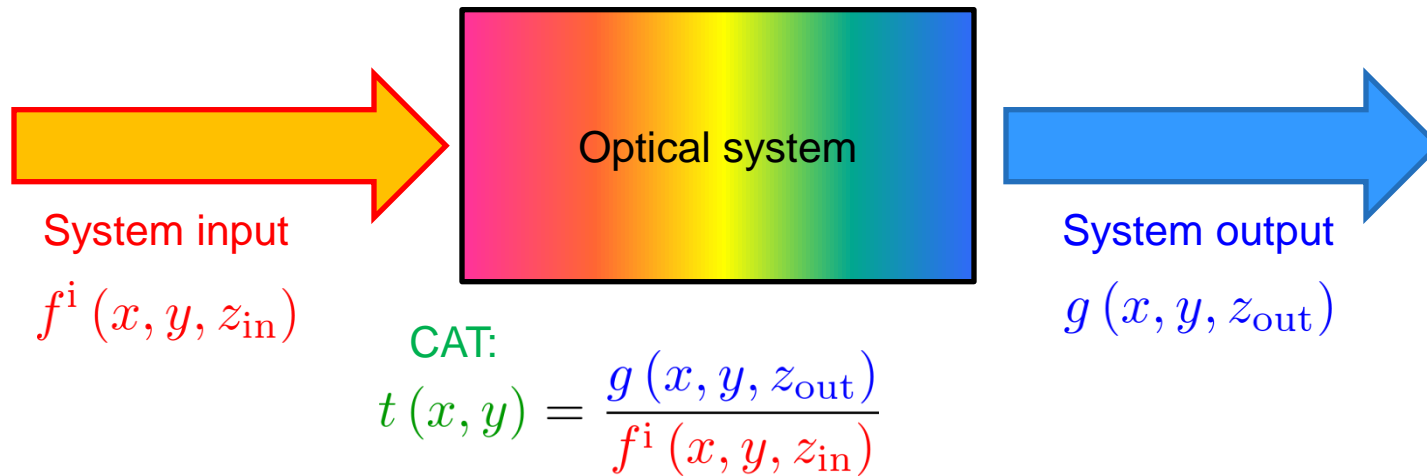
III. COMPARISON WITH OPTICS*

IV. EXAMPLES

❖ CONCLUSIONS & QUESTIONS

*J. W. Goodman, Introduction to Fourier Optics, 2nd edition, 1996.

The Complex Amplitude Transmittance (CAT)



Not limited to
bulky elements

Spatial Light
Modulators
(SLM)

EXAMPLE 4.1-1. Scanning. A thin transparency with complex amplitude transmittance $f(x, y) = \exp(j\pi x^2/\lambda f)$ introduces a phase shift $2\pi\phi(x, y)$ where $\phi(x, y) = -x^2/2\lambda f$, so that the wave is deflected at the position (x, y) by the angles $\theta_x = \sin^{-1}(\lambda\partial\phi/\partial x) = \sin^{-1}(-x/f)$ and $\theta_y = 0$. If $|x/f| \ll 1$, $\theta_x \approx -x/f$ and the deflection angle θ_x is directly proportional to the transverse distance x . This transparency may be used to deflect a narrow beam of light. If the transparency is moved at a uniform speed, the beam is deflected by a linearly increasing angle as illustrated in Fig. 4.1-6.

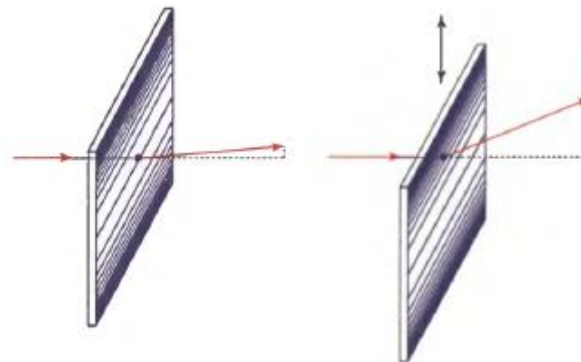


Figure 4.1-6 Using a frequency-modulated transparency to scan an optical beam.

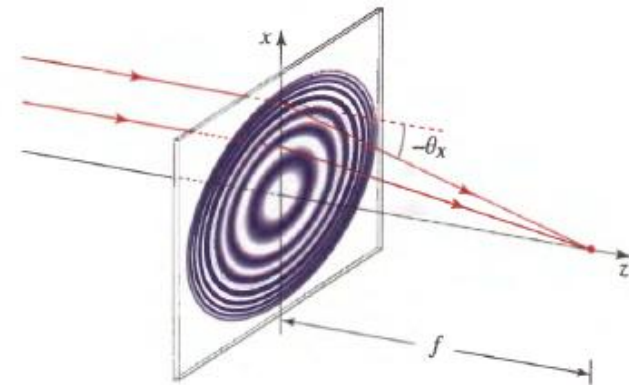
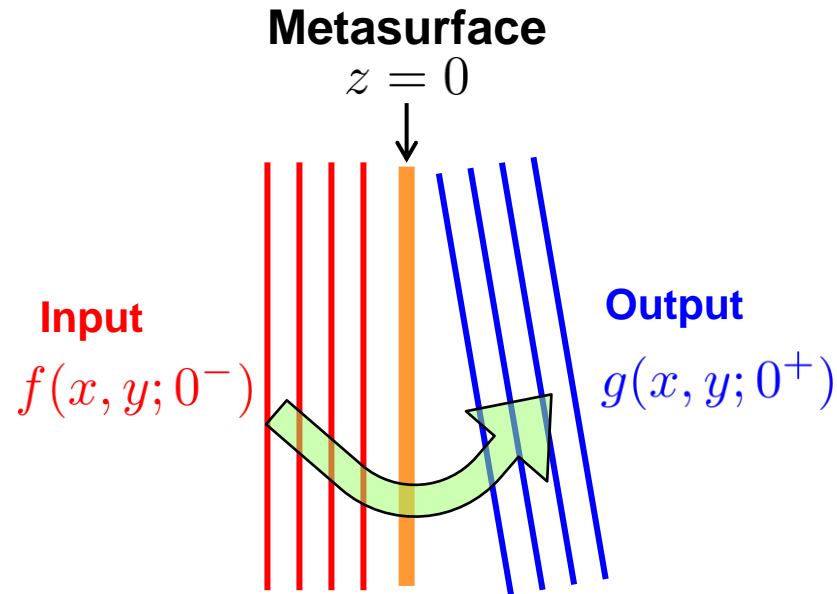


Figure 4.1-7 A transparency with transmittance $f(x, y) = \exp(j\pi x^2/\lambda f)$ bends the wave at position x by an angle $\theta_x \approx -x/f$ so that it acts as a cylindrical lens with focal length f .

B. E. A. Saleh and M. C. Teich,
Fundamentals of Photonics, 2007.



Momentum transformation (MT):

$$\eta(x, y) = \frac{g(x, y, 0^+)}{f(x, y, 0^-)}$$

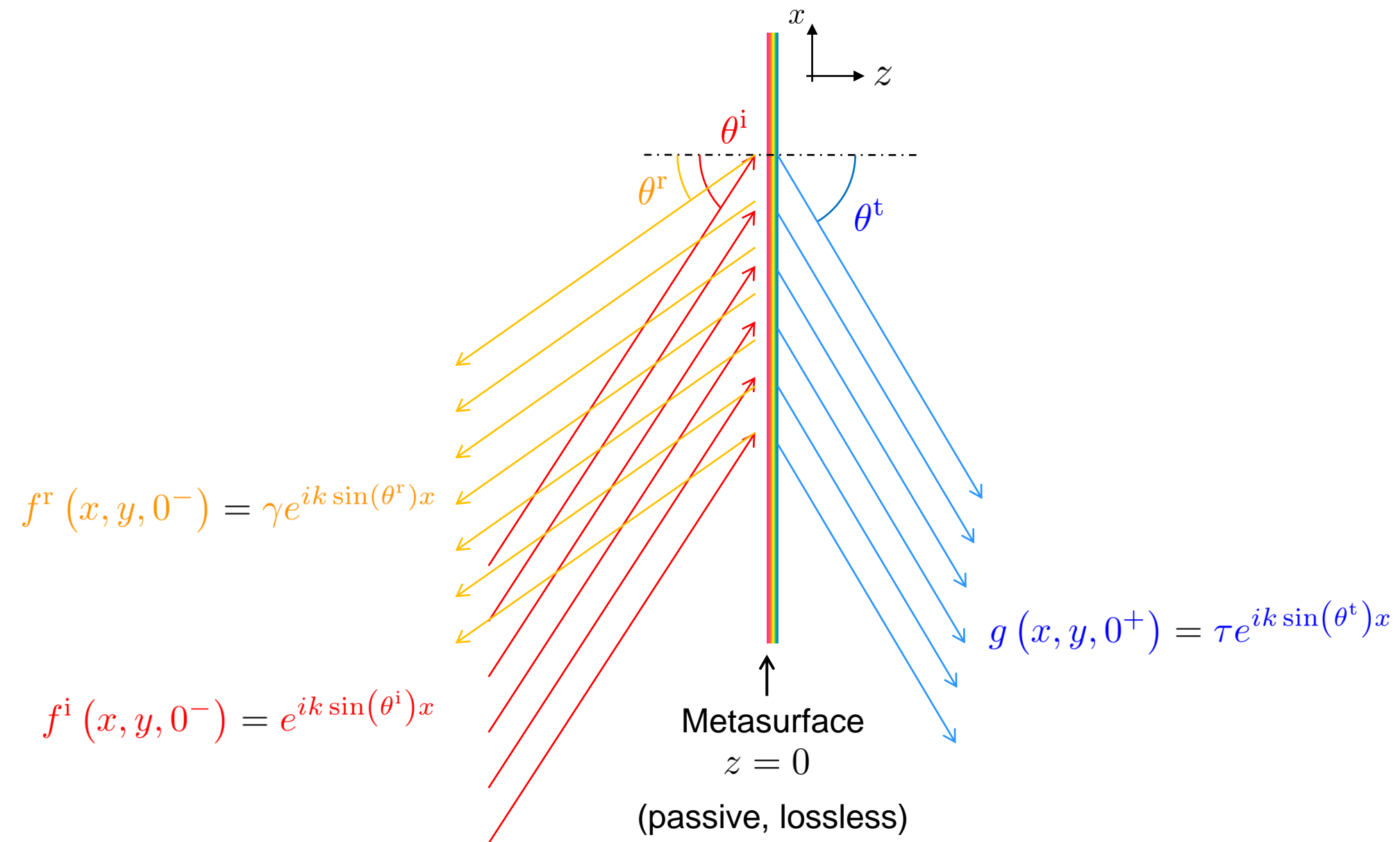
Complex amplitude transmittance (CAT):

$$t(x, y) = \frac{g(x, y, 0^+)}{f^i(x, y, 0^-)}$$

Definition of input:

$$f(x, y, 0^-) = f^i(x, y, 0^-) + f^r(x, y, 0^-)$$

Illustrative Example: Anomalous Reflection & Refraction



e.g. 2D problem, TE_y wave

CAT Approach

$$\begin{aligned} t(x, y) &= \frac{g(x, y, 0^+)}{f^i(x, y, 0^-)} \\ &= \frac{\tau e^{ik \sin(\theta^t)x}}{e^{ik \sin(\theta^i)x}} \\ &= \tau \exp(ikx[\sin(\theta^t) - \sin(\theta^i)]) \end{aligned}$$

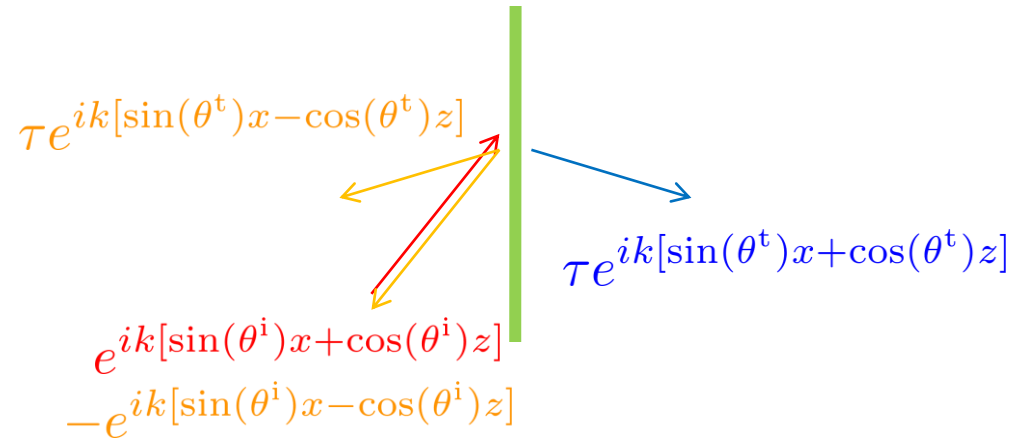
Boundary conditions

(2D-TE)

$$\begin{aligned} f(x, y, 0^-) &= g(x, y, 0^+) \\ \partial_z f(x, y, z)|_{z=0^-} &= \partial_z g(x, y, z)|_{z=0^+} \end{aligned}$$

$$f(x, y, z) = f^i(x, y, z) + f^r(x, y, z)$$

What about the reflected wave?



Cannot be used for synthesis:

Not the required behavior

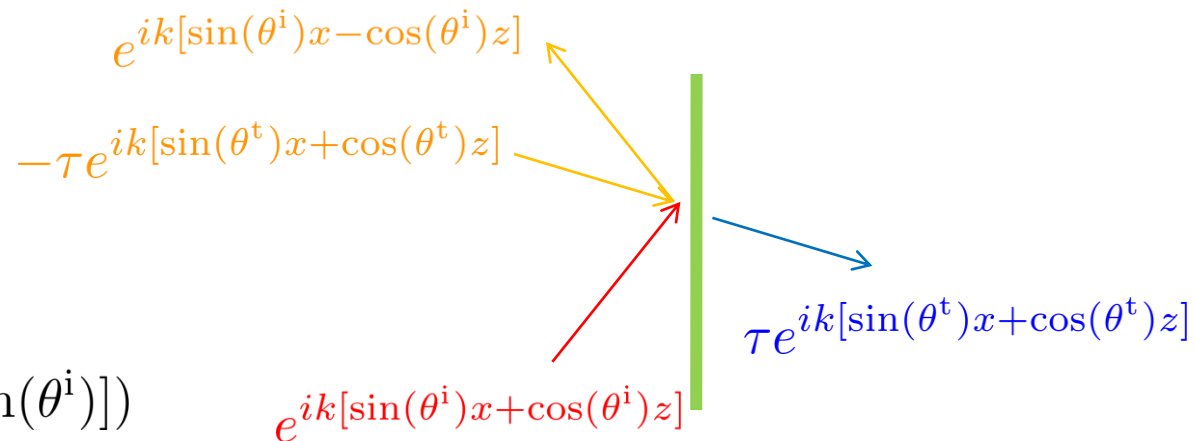
Violating passivity

$$\begin{aligned} f^r(x, y) &= g(x, y) - f^i(x, y) \\ &= \tau e^{ik \sin(\theta^t)x} - e^{ik \sin(\theta^i)x} \end{aligned}$$

$$\begin{aligned} f^r(x, y, z) &= \tau e^{ik[\sin(\theta^t)x - \cos(\theta^t)z]} \\ &\quad - e^{ik[\sin(\theta^i)x - \cos(\theta^i)z]} \end{aligned}$$

CAT Approach

$$\begin{aligned} t(x, y) &= \frac{g(x, y, 0^+)}{f^i(x, y, 0^-)} \\ &= \frac{\tau e^{ik \sin(\theta^t)x}}{e^{ik \sin(\theta^i)x}} \\ &= \tau \exp(ikx[\sin(\theta^t) - \sin(\theta^i)]) \end{aligned}$$



Energy Conservation

(same medium)

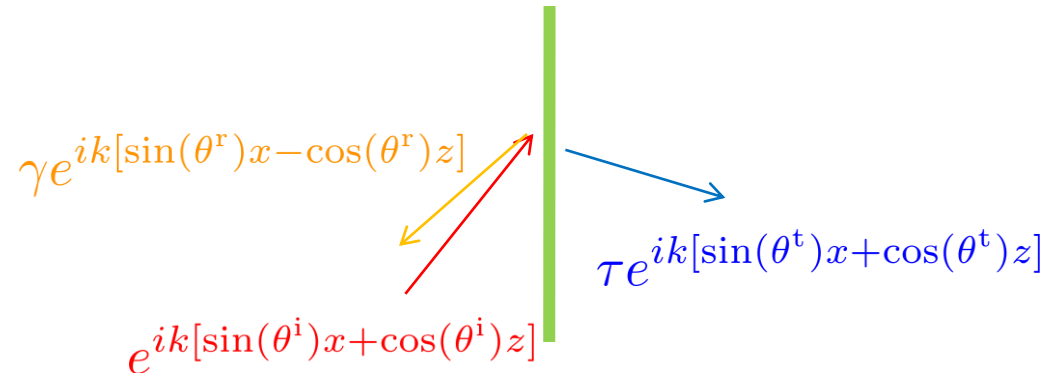
$$\begin{aligned} f^i(x, y, 0^-) &= f^r(x, y, 0^-) + g(x, y, 0^+) \\ &= f^r(x, y, 0^-) + t(x, y) f^i(x, y, 0^-) \end{aligned}$$

$$f^r(x, y, 0^-) = (1 - t(x, y)) f^i(x, y, 0^-)$$

Cannot be used for synthesis:
Not the required behavior
Violating causality

MT Approach

$$\begin{aligned}\eta(x, y) &= \frac{g(x, y, 0^+)}{f(x, y, 0^-)} \\ &= \frac{\tau e^{ik \sin(\theta^t)x}}{e^{ik \sin(\theta^i)x} + \gamma e^{ik \sin(\theta^r)x}}\end{aligned}$$



Boundary conditions

(in the sense of distributions)

$$\begin{aligned}g(x, y, 0^+) - f(x, y, 0^-) &= \Lambda(x, y) \\ \partial_z g(x, y, z)|_{z=0^+} - \partial_z f(x, y, z)|_{z=0^-} &= \underbrace{\Lambda'(x, y)}\end{aligned}$$

Can be used for synthesis:
Correct physical description

Related to $\eta(x, y)$

$$\Lambda(x, y) = (\eta(x, y) - 1) f(x, y, 0^-)$$

$$\begin{aligned}f^r(x, y, 0^-) &= g(x, y, 0^+) - f^i(x, y, 0^-) - \Lambda(x, y) \\ &= \gamma e^{ik \sin(\theta^r)x}\end{aligned}$$

Boundary conditions

(in the sense of distributions)

$$g(x, y, 0^+) - f(x, y, 0^-) = \Lambda(x, y)$$

$$\tilde{g}(k_x, k_y; 0^+) - \tilde{f}(k_x, k_y; 0^-) = \tilde{\Lambda}(k_x, k_y) \longleftrightarrow \tilde{g}(k_x, k_y; 0^+) - \tilde{f}(k_x, k_y; 0^-) = 0$$

Paradox?

Momentum conservation (physics)

Discontinuous momenta

(require distribution notation)

$$\langle \tilde{g}, \tilde{\phi} \rangle - \langle \tilde{f}, \tilde{\phi} \rangle = \langle \tilde{\Lambda}, \tilde{\phi} \rangle$$

$$\langle \tilde{\Lambda}, \tilde{\phi} \rangle = \langle 0, \tilde{\phi} \rangle$$

Particular solution:

$$\tilde{\Lambda} = 0$$

Homogeneous solution:

$$\text{supp} [\tilde{\Lambda}] = 0 \longrightarrow \tilde{\Lambda} = \Lambda(x, y) \delta(k_x, k_y)$$

From momentum transformation: $\Lambda = (\eta - 1) f$
(Surface response function)

M. A. Salem and C. Caloz, Opt. Express, 22, 2014.

Straightforward generalization to vector case !

I. MOTIVATION

II. MT METHOD

III. COMPARISON WITH OPTICS

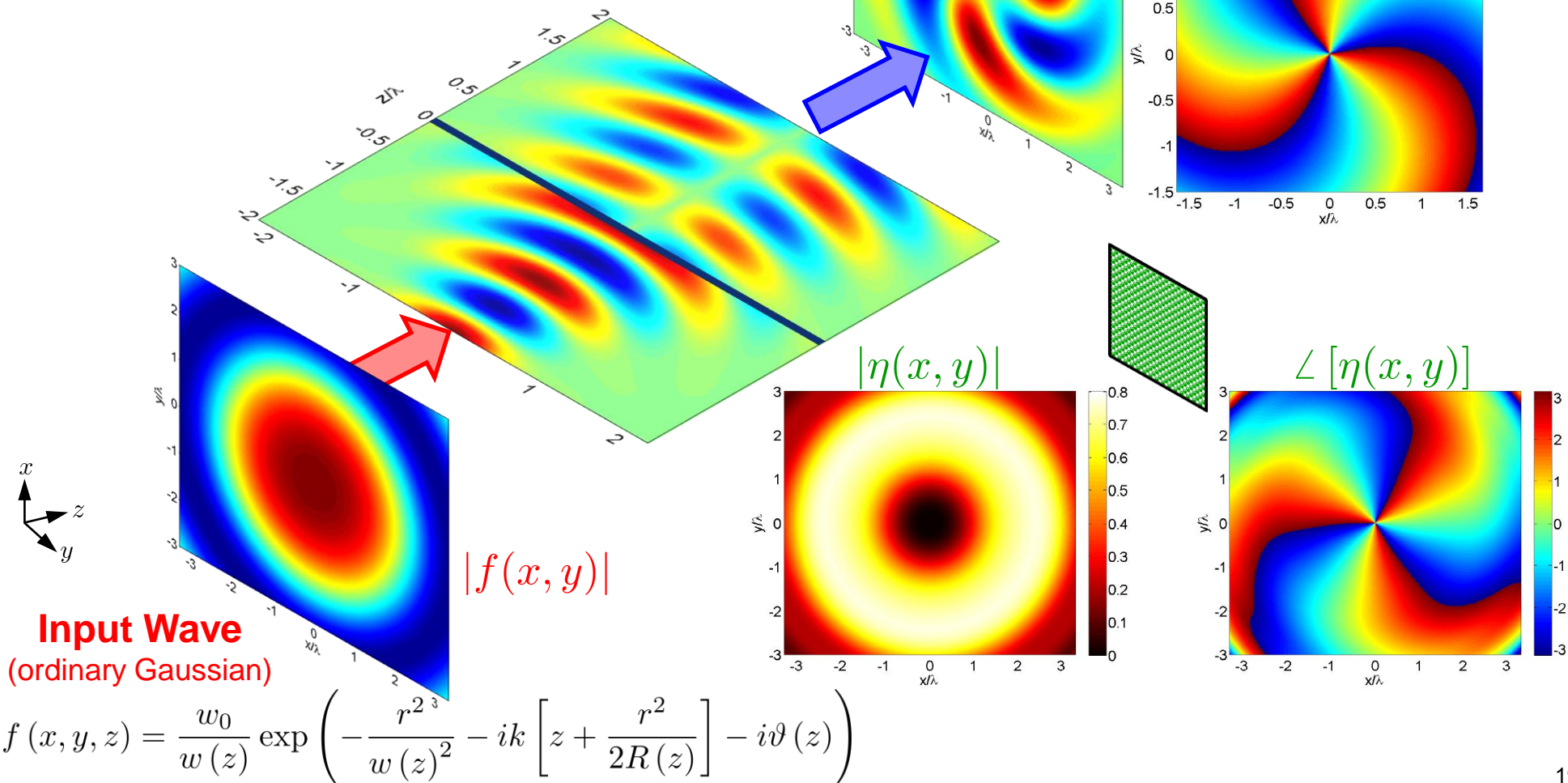
IV. EXAMPLES

❖ CONCLUSIONS & QUESTIONS

Output Wave

(vortex hypergeometric Gaussian)

$$g(x, y, z) = \frac{i^{|m|+1} \xi^{p/2}}{[\xi + i]^{1+|m|+p/2}} \left[\frac{\rho}{w_0} \right]^{|m|} \exp \left(\frac{-i\rho^2}{w_0^2 [\xi + i]} + im\phi \right) \\ \times \frac{\Gamma(1 + |m| + \frac{p}{2})}{\Gamma(|m| + 1)} {}_1F_1 \left(-\frac{p}{2}, |m| + 1; \frac{\rho^2}{w_0^2 \xi [\xi - i]} \right)$$

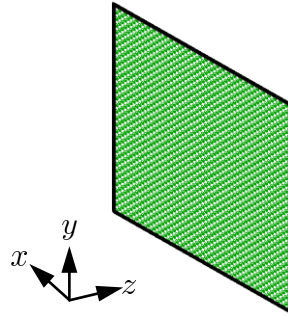


Delay-Start Airy Beam Generation

Input Wave

(normally incident plane wave)

$$f(z) = e^{ikz}$$

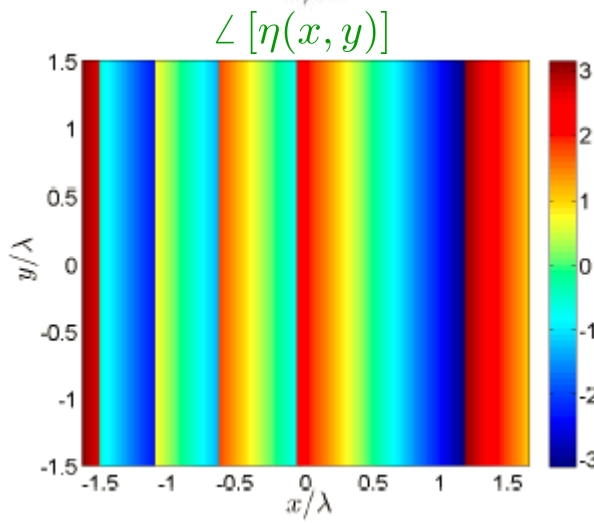
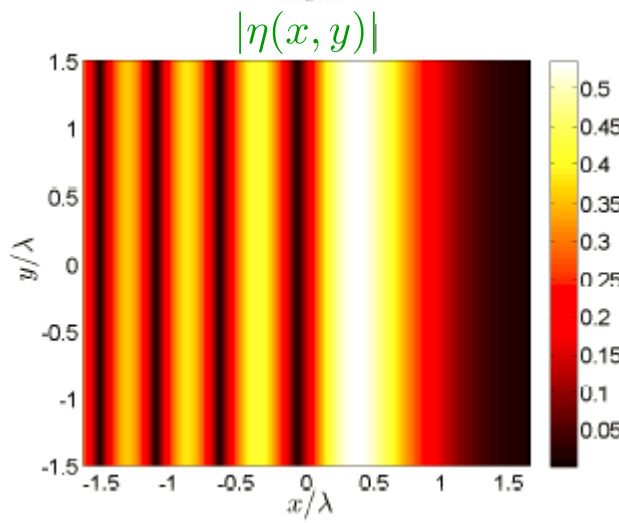
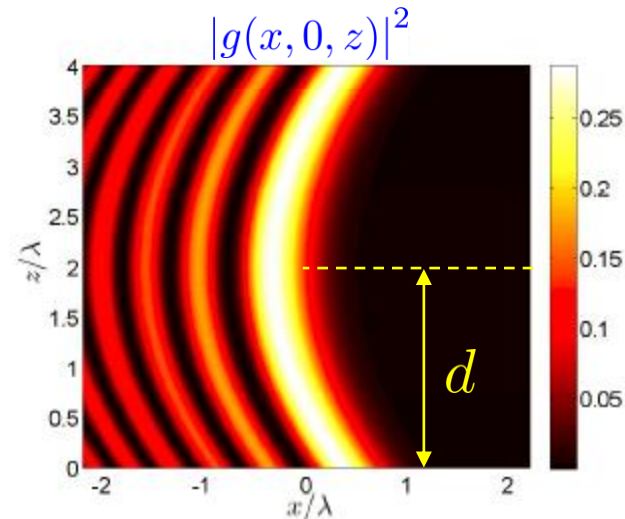
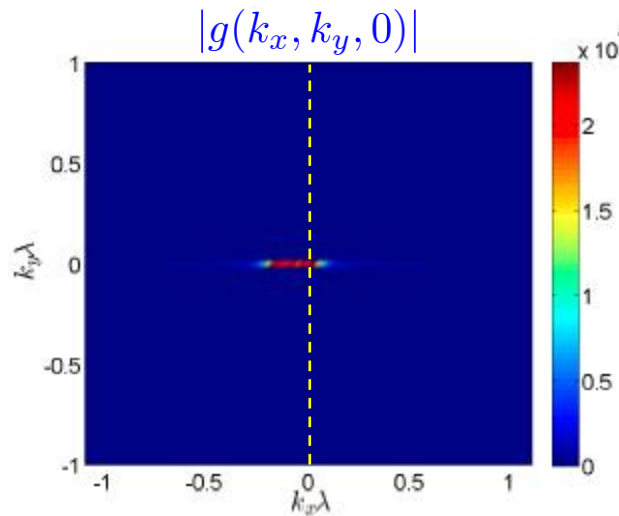


Output Wave (Airy beam with start at $z = d$)

$$g(x, z) = \text{Ai} \left(s - \frac{\zeta^2}{4} \right) \exp \left(i \left[\frac{s\zeta}{2} - \frac{\zeta^3}{12} \right] \right)$$

$$\tilde{\varphi}(k_x; k, d) = \exp \left(-i \sqrt{k^2 - k_x^2} d \right)$$

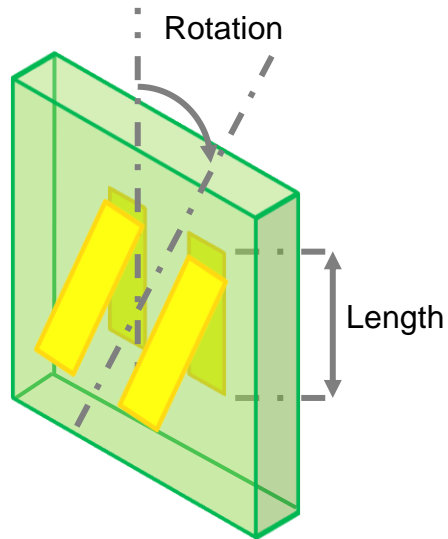
$$s = x/x_0, \quad \zeta = z/kx_0^2$$



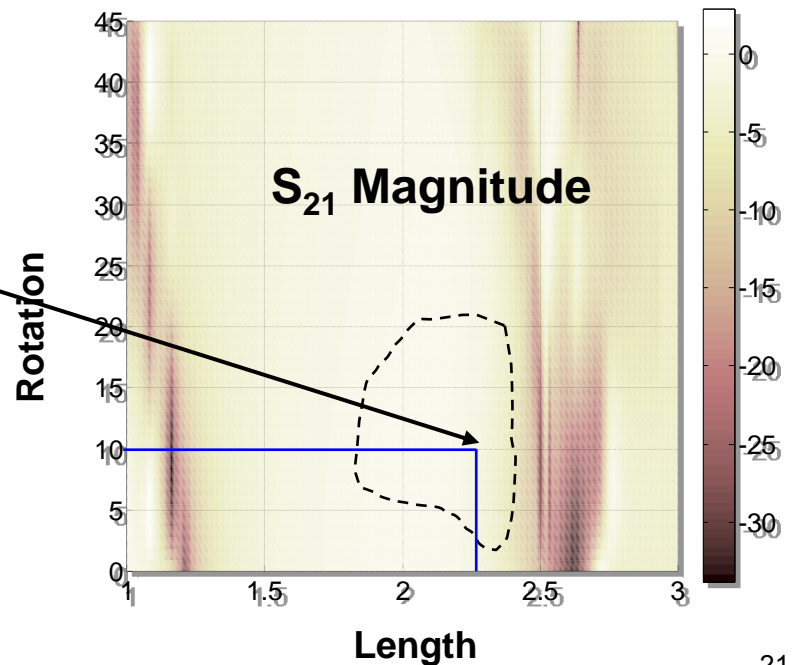
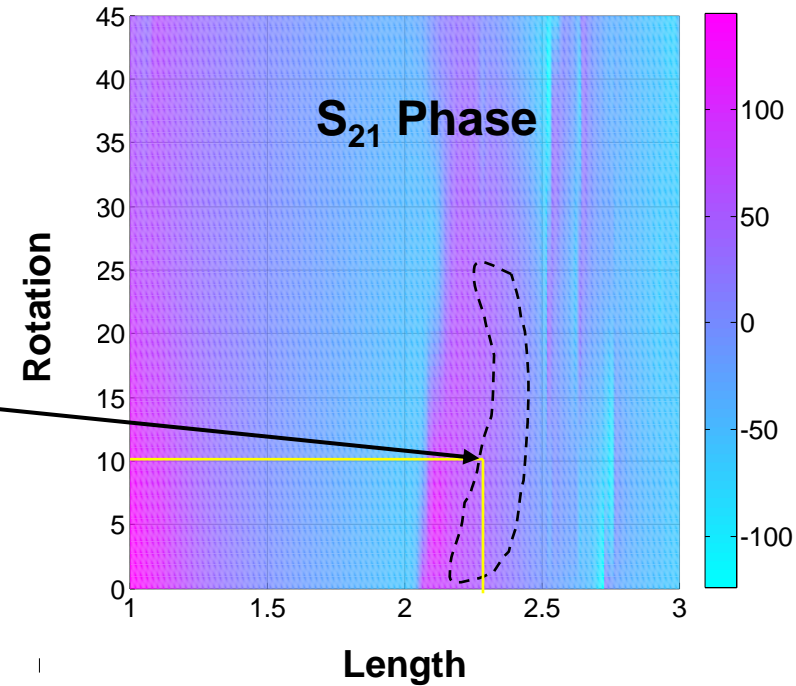
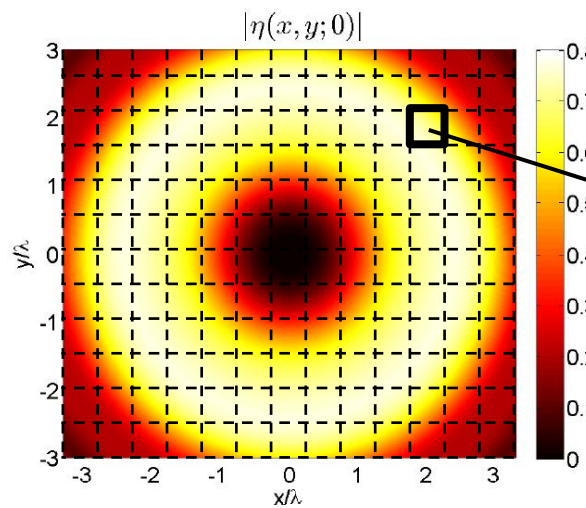
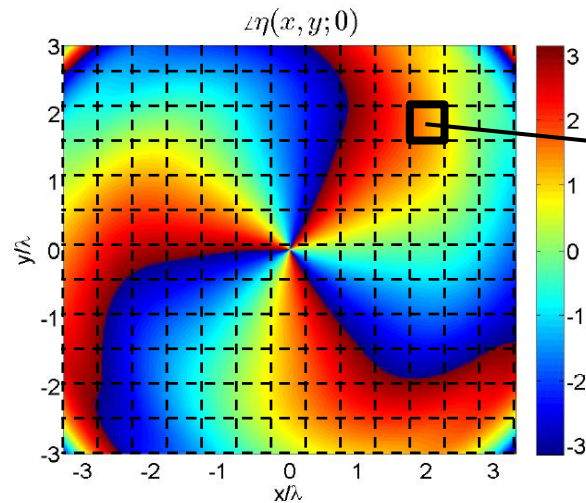
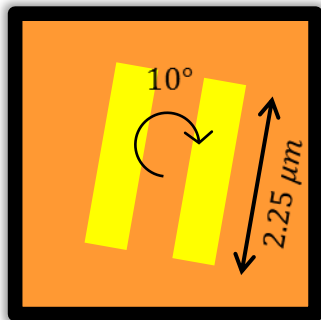
Implementation via Parametric Mapping

$$\eta(x, y) \Rightarrow \begin{bmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{bmatrix} \text{ (de-embedded)}$$

Scatterer topology



Pixel Dimensions



I. MOTIVATION

II. MT METHOD

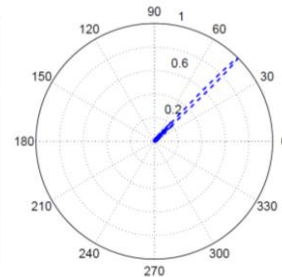
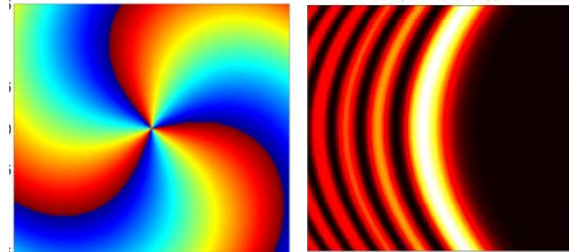
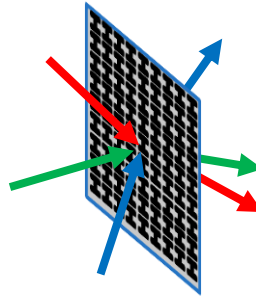
III. COMPARISON WITH OPTICS

IV. IMPLEMENTATION SUGGESTION

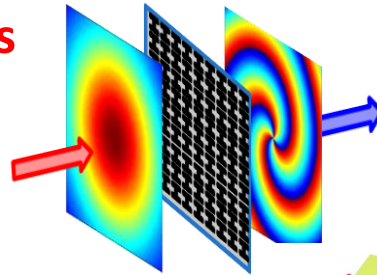
❖ CONCLUSIONS & QUESTIONS

Conclusions

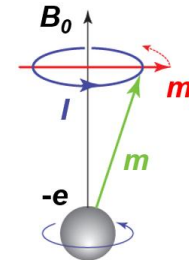
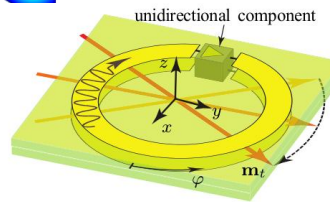
- Metasurfaces:
unprecedented EM control



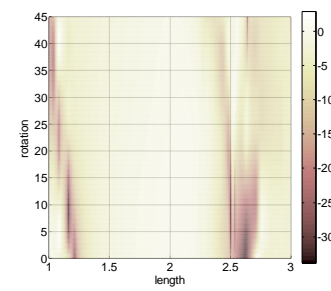
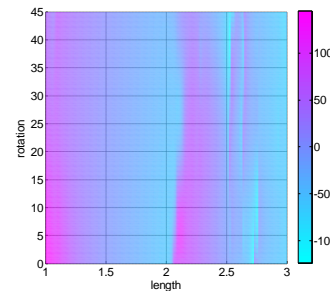
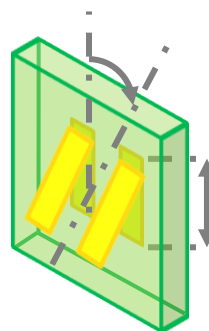
- No universal synthesis
– until now



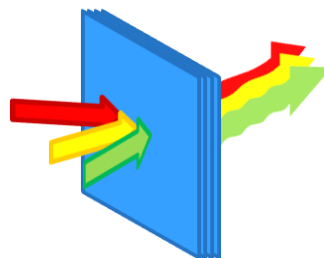
- Synthesis based on
transformation in momentum
space for physical insight



- Implementation using
lookup maps



- Potential for myriads
of applications



K. Achouri, M. A. Salem and C. Caloz,
“General Metasurface Synthesis Based on
Susceptibility Tensors,” arXiv:1408.0273.